# Theoretical characterization of pentazole anion with metal counter ions. Calculated and experimental ${ }^{15} \mathrm{~N}$ shifts of aryldiazonium, -azide and -pentazole systems $\dagger$ 

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Theoretical studies of proposed structures for $\mathrm{NaN}_{5}, \mathrm{KN}_{5}, \mathrm{Mg}\left(\mathrm{N}_{5}\right)_{2}, \mathrm{Ca}\left(\mathrm{N}_{5}\right)_{2}$, and $\mathrm{Zn}\left(\mathrm{N}_{5}\right)_{2}$ metal complexed pentazole anions have been carried out with the RHF, MP2, MCSCF, and DFT theoretical methods. Additional DFT calculations were performed on $\mathrm{MgN}_{5} \mathrm{Cl}, \mathrm{CaN}_{5} \mathrm{Cl}$, and $\mathrm{ZnN}_{5} \mathrm{Cl}$ pentazoles. The structures considered are unidentate I, bidentate II, and metallocene-like III. For $\mathrm{Mg}, \mathrm{Na}, \mathrm{K}$, and Ca pentazoles at every level of theory, II is the most energetically favoured, followed by I, then III. Complex I is preferred with Zn complexes due to favourable d orbital interactions. For double ring complexes only $\mathbf{I I}(\mathbf{I}$ for Zn ) with perpendicular rings has all positive vibrational frequencies. For single ring complexes, both II (I for Zn ) and III have all positive vibrations. Structure $\mathbf{I}(\mathbf{I I}$ for $\mathbf{Z n})$ is a transition state structure for metal ion rotation around the ring $\left(E_{\mathrm{a}} 5-10 \mathrm{kcal} \mathrm{mol}^{-1}\right)$. N atom chemical shifts relative to $\mathrm{NH}_{3}$ and nitromethane were calculated for each species using the lowest energy configuration and the B3LYP//6-311++G(2d,p) method on the B3LYP//6-31G(d) optimised geometry. Additional calculations were done for 1 -arylpentazoles, 1 -arylpentazene, aryl azides, and aryldiazonium ions. Calculated ${ }^{15} \mathrm{~N}$ NMR shifts were within 20 ppm of experiment. Time dependent B3LYP/6-31G(d) and B3LYP/6-311+G(d) calculations were performed on all stable species. All ${ }^{1}(\pi, \pi)$ transitions were calculated to be below 180 nm , while the ${ }^{1}(\mathrm{n}, \pi)$ transitions were below 210 nm . The lowest energy transitions are from the lone pairs to the empty metal s orbital. For Mg and Zn these transitions are at $\sim 220 \mathrm{~nm}$. For $\mathrm{Na}, \mathrm{Ca}$ and K the transitions are considerably lower in energy, $\sim 250 \mathrm{~nm}$.

## Introduction

Substituted 1-arylpentazoles have been synthesised and characterised for over forty years. ${ }^{1}$ There have been no reports yet of the preparation of the parent pentazole $\mathrm{HN}_{5}$ or its metal salts. We report here the results of theoretical calculations on pentazole anion and various metal cation-pentazole anion complexes in terms of their structures, vibrational frequencies, electronic transitions, and ${ }^{15} \mathrm{~N}$ chemical shifts. Metallocene-like structures have been proposed ${ }^{2}$ but no studies have appeared where vibrational frequencies have been calculated in order to characterise the structures as first or higher order transition states. Yet another structure (Scheme 1) could be expected based on the known aryl substituted pentazole structure. ${ }^{1-3}$


Scheme 1
In the case of $\mathrm{Na}^{+}$and $\mathrm{K}^{+}$, there is only substitution to one ring. For the divalent $\mathrm{Mg}^{2+}, \mathrm{Ca}^{2+}$, and $\mathrm{Zn}^{2+}$ ions there is possible substitution by two rings or by one ring and another anion. For the purposes of this study, we have chosen the former for an analysis by different levels of theory and the latter only by Density Functional Theory (DFT). In exploring the
$\dagger$ Electronic supplementary information (ESI) available: Table S1selected bond lengths calculated using the $6-31 \mathrm{G}(\mathrm{d})$ basis set with various theoretical methods. See http://www.rsc.org/suppdata/p2/b1/ b103852b/

la

lb


Ila

lib


IIla


IIIb

Scheme 2 N at all corners.
hypersurfaces, three principal structures were found (Scheme 2, N at all corners). They are $\mathbf{I}$, with one point of attachment to the ring (unidentate), II two points (bidentate), and III centrally over the ring (metallocene-like). In the double rings, the rings can rotate from $\left(D_{2 \mathrm{~d}}\right)$ perpendicular (a) to $\left(D_{2 \mathrm{~h}}\right)$ planar (b) in I and II, and in structure III from $\left(D_{5 \mathrm{~h}}\right)$ staggered (a) to $\left(D_{5 \mathrm{~h}}\right)$ eclipsed (b).

## Computational details

A number of computational methods incorporated into the Gaussian98A7 series of programs were used in this study. ${ }^{4}$ All geometry optimisations were carried out at the restricted Hartree-Fock (RHF), RMP2, and RB3LYP ${ }^{5}$ levels for $\mathrm{NaN}_{5}$,
$\mathrm{KN}_{5}, \mathrm{Mg}\left(\mathrm{N}_{5}\right)_{2}, \mathrm{Ca}\left(\mathrm{N}_{5}\right)_{2}$, and $\mathrm{Zn}\left(\mathrm{N}_{5}\right)_{2}$. Results from both the standard split valence plus polarisation $6-31 \mathrm{G}(\mathrm{d})$ basis set and the larger $6-311+G(d)$ set containing diffuse functions were obtained. Additionally, the CASSCF method, which included the four highest and four lowest $\pi$ molecular orbitals (MO), was employed for $\operatorname{Mg}\left(\mathrm{N}_{5}\right)_{2}$ and calculations using the two highest and two lowest $\pi$ MOs were performed for $\mathrm{NaN}_{5}$. As $\mathrm{NaN}_{5}$ contains one ring and $\operatorname{Mg}\left(\mathrm{N}_{5}\right)_{2}$ two, these can be considered to be nearly equivalent level calculations. Stability of the restricted wavefunctions ${ }^{6}$ was verified for all the Na and Mg structures. Normal mode analysis was performed to ascertain the nature of structures identified as stationary points on the potential energy surfaces. Vibrational frequencies reported here were not empirically scaled. For the divalent cations, RB3LYP//6-31G(d) and $6-311+G(d)$ calculations were carried out for the cases where one of the rings was replaced by a Cl atom placed in a linear ring-M-Cl conformation. Figures with orbitals were generated with the Molekel Program. ${ }^{7}$

The N atom chemical shifts were calculated according to the Gauge-Independent Atomic Orbital (GIAO) method. ${ }^{8}$ RB3LYP//6-311++G(2d,p) level calculations were carried out on RB3LYP//6-31G(d) optimised structures. Two test cases for $\mathrm{NaN}_{5}$ were performed to evaluate the effect of geometry change on the calculated chemical shifts by using the RB3LYP// $6-311+G(d)$ optimised structures and by using conformer III. Nitromethane was used as the reference molecule and extra calculations on 1 - and 2 -substituted methyltetrazoles were done as a test of the theoretical method compared to the experimental. ${ }^{8}$ Using the values for 24 N atoms in ten different azoles Witanowski et al. ${ }^{9}$ found a linear relationship between the experimental results and that for Coupled Hartree-Fock (CHF) $/ 6-31++\mathrm{G}(\mathrm{d}, \mathrm{p})$ of $\sigma_{\text {exp }}=\left(0.8804 \sigma_{\text {calc }}+112.56\right)$. Finally, a time dependent DFT method ${ }^{10}$ (TD-DFT) was used to calculate the electronic transitions of $\mathrm{MgN}_{5} \mathrm{Cl}, \mathrm{NaN}_{5}, \mathrm{KN}_{5}, \mathrm{CaN}_{5} \mathrm{Cl}$ and $\mathrm{ZnN}_{5} \mathrm{Cl}$ at the $6-311+\mathrm{G}(\mathrm{d})$ levels.

## Results and discussions

In all cases, structures II or IIa were calculated to be the lowest in energy and have no negative vibrational frequency, except for $\mathrm{Zn}^{2+}$ where $\mathbf{I}$ or $\mathbf{I a}$ are lowest and have no negative frequency. The relative energies in $\mathrm{kcal} \mathrm{mol}^{-1}$ are reported in Table 1. The output files for all calculations in this study can be found at: http://camchem.rutgers.edu/~burke.

In all cases, the b structures have one more negative frequency than the a structures. Vector analysis (Scheme 3)

shows each $\mathbf{b}$ structure to contain a rotational transition state (TS) around the cation. The negative frequency in the single ring structure $\mathbf{I}$ is an in-plane wag of the $\mathrm{Na}^{+}$or $\mathrm{K}^{+}$cation between two equivalent II structures. The two negative frequencies in the $\mathrm{Mg}^{2+}$ and $\mathrm{Ca}^{2+} \mathbf{I a}$ structures (IIa in $\mathrm{Zn}^{2+}$ ) are due to the coupled symmetric and asymmetric combinations of the wag modes found in I. The three modes in Ib include the two for Ia plus the ring rotation mode also found in IIb The four negative frequencies in IIIa correspond to two pairs of symmetry-coupled modes, the first pair for slippage of the cation between the rings and the other for movement of the


Fig. 1 Plots of eight DFT MO are given for the $N_{5}$ anion and the $\mathrm{NaN}_{5}$ ion pair. Both sets of eight HOMOs include the five lone pair orbitals and the three $p$ orbitals. The LUMOs of the anion consist of a degenerate pair of p MOs; while the LUMO in all ion pair cases is mainly of s character with the fourth and fifth p MO dispersed among the metals' virtual p and d orbitals. The highest shown anion MO is a low lying Rydberg orbital. $D_{5 \mathrm{~h}}$ and $C_{2 \mathrm{v}}$ group symmetries are also given for the MOs. $6-311+\mathrm{G}(\mathrm{d})$ basis set.
rings toward structure IIa. The fifth negative frequency in IIIb is due to ring rotation from the eclipsed to staggered conformations. There are no negative frequencies for conformer III with $\mathrm{Na}^{+}$or $\mathrm{K}^{+}$.

In the $(\mathrm{M}-\mathrm{Cl})^{+}$cases $(\mathrm{M}=\mathrm{Mg}, \mathrm{Ca})$, the single negative frequency for $\mathbf{I}$ also corresponds to a TS between equivalent II structures. With $\mathbf{M}=\mathbf{Z n}, \mathbf{I I}$ is a TS between two equivalent $\mathbf{I}$ structures. Although the III structure for $\mathrm{Na}, \mathrm{K}$, and $\mathrm{M}-\mathrm{Cl}$ is an equilibrium structure with all positive frequencies, it is higher in energy than II or I in all cases. The relative energies range from a few $\mathrm{kcal} \mathrm{mol}{ }^{-1}$ for K and $\mathrm{Na}, c a .10$ for $\mathrm{Ca}, c a .20$ for Mg , and to over 30 for Zn . A stationary state for the III conformation was not found for $\mathrm{ZnClN}_{5}$ using the $6-31 \mathrm{G}(\mathrm{d})$ basis set.
Inspection of Table 1 shows little effect on relative energies by differences in theoretical method used but a rather general effect by change in basis set. The DFT results for either basis set lie between those found with the RHF and RMP2 methods. Table S1 (provided as supplementary material $\dagger$ ) gives the optimised bond lengths in $\AA$ found with the various theoretical methods and basis sets. In all methods, the presence of the cation on the side of the ring induces alternating bond lengths, thus less aromatic character
Fig. 1 contains plots of the MOs in the anion and $\mathrm{NaN}_{5}$ calculated with the B3LYP//6-311+(d) method. The eight highest occupied molecular orbitals of the anion consist of five lone pair orbitals, two pairs of which are degenerate, and three $\pi$ orbitals of which one pair is degenerate. The $D_{5 h}$ symmetry of the anion is reduced to $C_{2 v}$ in presence of the cation in either the I or II conformation. The three lowest unoccupied molecular orbitals in the anion consist of the highest degenerate $\pi$ orbitals and the $\mathrm{A}_{1}{ }^{\prime} \sigma^{*}$ MO which resembles the lowest occupied $\mathrm{A}_{1}{ }^{\prime}$ lone pair MO but consisting of Rydberg 3p AOs with higher radial nodes (higher regions not shown on Fig. 1). The LUMO of all ion pair structures consists of the metal vacant s orbital. The vacant metal $\mathrm{p}_{x}, \mathrm{p}_{y}$, and $\mathrm{p}_{z}$ and $\mathrm{d} A O$ levels are interspersed with the fourth and fifth $\pi$ orbitals according to the various metal cations. However, in Zn the d AOs are occupied.
The presence of the Cl atom introduces three occupied p -like orbitals making the top 11 orbitals significant to bonding. Fig. 2 shows the 11 MOs for the Zn case for both the I and II conformations, plus three of the virtual MOs. Each conformation gives a remarkably similar pattern of orbitals among

Table 1 Relative total energies ( $\mathrm{kcal} \mathrm{mol}^{-1}$ ) for metal-pentazole complexes calculated by theoretical methods ( $6-31 \mathrm{G}^{*}$ prime on method indicates the $6-311+G^{*}$ basis set)

| Metal | Method | Ia | Ib | IIa | IIb | IIIa | IIIb |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Mg | RHF | 9.19 | 9.29 | 0 | 3.90 | 41.88 | 41.92 |
|  | RHF ${ }^{\prime}$ | 7.80 | 7.96 | 0 | 1.04 | 44.27 |  |
|  | MP2 | 15.36 | 15.49 | 0 | 1.30 | 35.54 | 35.54 |
|  | MP2 ${ }^{\prime}$ | 12.57 | 12.92 | 0 | 1.59 | 39.58 | 39.77 |
|  | CAS(88) | 4.80 | 4.91 | 0 |  | $70^{a}$ |  |
|  | B3LYP | 11.59 | 11.70 | 0 | 1.13 | 41.10 | 41.14 |
|  | B3LYP ${ }^{\prime}$ | 9.60 | 9.77 | 0 | 1.15 | 47.59 | 47.66 |
|  | B3LYPCl | 4.74 | - | 0 | - | 19.18 | - |
|  | B3LYPCl ${ }^{\prime}$ | 3.81 | - | 0 | - | 22.68 | - |
| Na | RHF | 5.42 |  | 0 |  | 5.21 |  |
|  | RHF ${ }^{\prime}$ | 4.56 |  | 0 |  | 7.04 |  |
|  | MP2 | 6.26 |  | 0 |  | 5.30 |  |
|  | MP2 ${ }^{\prime}$ | 4.57 |  | 0 |  | 6.09 |  |
|  | CAS(44) | 4.87 |  | 0 |  | 12.46 |  |
|  | CAS(44) ${ }^{\prime}$ | 4.03 |  | 0 |  | 14.53 |  |
|  | B3LYP | 6.11 |  | 0 |  | 5.12 |  |
|  | B3LYP ${ }^{\prime}$ | 4.59 |  | 0 |  | 7.66 |  |
| K | RHF | 5.45 |  | 0 |  | 2.28 |  |
|  | RHF ${ }^{\prime}$ | 4.79 |  | 0 |  | 3.32 |  |
|  | MP2 | 6.75 |  | 0 |  | 0.87 |  |
|  | MP2 ${ }^{\prime}$ | 5.49 |  | 0 |  | 1.55 |  |
|  | B3LYP | 5.99 |  | 0 |  | 2.11 |  |
|  | B3LYP ${ }^{\prime}$ | 4.98 |  | 0 |  | 3.74 |  |
| Ca |  |  |  |  | 3.90 |  |  |
|  | RHF ${ }^{\prime}$ | 18.62 |  | 0 |  | 25.92 | 25.94 |
|  | MP2 | 23.27 |  | 0 | $0.29$ | 16.87 | 16.88 |
|  | B3LYP | 21.61 | 21.72 | 0 | 0.45 | 20.52 | 20.51 |
|  | B3LYP ${ }^{\prime}$ | 21.05 | 21.24 | 0 | 1.08 | 26.31 | 26.32 |
|  | B3LYPCl | 10.52 | - | 0 | - | 9.84 | - |
|  | B3LYPCl ${ }^{\prime}$ | 10.03 | - | 0 | - | 11.36 | - |
| Zn | RHF | 0 | 0.40 | 14.17 | 16.12 | 68.55 | 68.60 |
|  | RHF ${ }^{\prime}$ | 0 |  | 17.12 |  | 71.39 | 71.45 |
|  | MP2 | 0 |  | 13.57 | 16.97 | 73.68 | 73.88 |
|  | MP2 ${ }^{\prime}$ | 0 |  | 14.07 |  | 67.40 | 67.33 |
|  | B3LYP | 0 | 0.52 | 12.12 | 16.66 | 69.39 | 69.43 |
|  | B3LYP' ${ }^{\text {B }}$ | 0 | 0.28 | 17.47 | 19.60 | ${ }^{\text {a }}$ | 81.74 |
|  | B3LYPCl | 0 | - | 7.23 | - | 39 | - |
|  | B3LYPCl ${ }^{\prime}$ | 0 | - | 9.88 | - | 39.22 | - |

${ }^{a}$ Stationary state not found, opens towards I or II.


Fig. 2 The eleven DFT HOMOs and three of the unoccupied DFT MOs are given for the $\mathbf{I}$ conformation of $\mathrm{ZnClN}_{5}$ (two left columns) and the II conformation (two right columns). $C_{2 v}$ group assignments are made for both. $6-311+\mathrm{G}(\mathrm{d})$ basis set.
all the cations. For example, the splitting of the $\mathrm{B}_{1} \mathrm{MOs}$ (third and seventh) in Zn (conformation I) is only slightly larger than that for Mg (conformation $\mathbf{I}), 0.0125$ to 0.0110 au , respectively.

The major reason for the lower energy of the I conformation compared to II in the Zn case might be due to the combination


Fig. 3 Energies (au) and plots for the 6-311+G(d) DFT orbitals for $\mathrm{ZnClN}_{5}$, conformation I in left column, II in right. Bottom five are the MOs with major d AO components; the top orbital in each column is displayed in Fig. 2.
of the occupied d orbitals with the lone pair or $\pi$ orbitals. Fig. 3 gives the five $d$ orbitals plus the combination with the lowest lone pair orbital in the I conformation and the five $d$ orbitals with the lowest $\pi$ orbital in the II conformation. There is significantly more splitting with the $\mathrm{d}_{z^{2}}$ orbital and the lone pair in the I conformation than with a $\mathrm{d}_{x y}$ orbital and the lowest $\pi$ orbital in the II conformation.

Table 2 Vibrational frequencies $\left(\mathrm{cm}^{-1}\right)$ calculated with the B3LYP/6-311+G(d) method for the lowest energy configurations. The $C_{2 v}$ symmetry designations of the metal complexed pentazoles are derivatives of those of the $D_{5 \mathrm{~h}}$ pentazole anion (in parentheses)

| Molecule | $\mathrm{N}_{5}{ }^{-}$ | $\mathrm{N}_{5}{ }^{-a}$ | $\mathrm{NaN}_{5}$ | $\mathrm{KN}_{5}$ | $\mathrm{MgN}_{5} \mathrm{Cl}$ | $\mathrm{CaN}_{5} \mathrm{Cl}$ | $\mathrm{ZnN}_{5} \mathrm{Cl}$ |
| :---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\mathrm{B}_{2}\left(\mathrm{E}_{1}{ }^{\prime}\right)$ | 1239.2 | 1286.1 | 1289.2 | 1281.5 | 1305.2 | 1299.6 | 1350.3 |
| $\mathrm{~A}_{1}\left(\mathrm{E}_{1}{ }^{\prime}\right)$ | 1239.2 | 1286.1 | 1230.5 | 1232.4 | 1248.7 | 1239.8 | 1275.0 |
| $\mathrm{~A}_{1}\left(\mathrm{~A}_{1}{ }^{\prime}\right)$ | 1182.0 | 1222.4 | 1216.5 | 1210.4 | 1217.2 | 1222.1 | 1218.0 |
| $\mathrm{~B}_{2}\left(\mathrm{E}_{2}{ }^{\prime}\right)$ | 1114.3 | 1124.2 | 1123.1 | 1123.0 | 1125.3 | 1123.1 | 1106.0 |
| $\mathrm{~A}_{1}\left(\mathrm{E}_{2}{ }^{\prime}\right)$ | 1114.3 | 1124.2 | 1111.8 | 1113.0 | 1104.3 | 1108.1 | 1130.2 |
| $\mathrm{~A}_{1}\left(\mathrm{E}_{2}{ }^{\prime}\right)$ | 995.8 | 1059.5 | 1025.3 | 1025.3 | 1036.5 | 1011.3 | 908.0 |
| $\mathrm{~B}_{2}\left(\mathrm{E}_{2}{ }^{\prime}\right)$ | 995.8 | 1059.5 | 980.3 | 988.4 | 953.4 | 973.0 | 1026.5 |
| $\mathrm{~B}_{1}\left(\mathrm{E}_{2}{ }^{\prime \prime}\right)$ | 767.2 | 782.7 | 766.2 | 765.7 | 756.4 | 756.7 | 731.8 |
| $\mathrm{~A}_{2}\left(\mathrm{E}_{2}{ }^{\prime \prime}\right)$ | 767.2 | 782.7 | 764.0 | 762.2 | 758.3 | 754.6 | 759.5 |
| $\mathrm{~A}_{1}$ |  |  | 308.7 | 230.0 | 584.1 | 412.6 | 483.8 |
| $\mathrm{~B}_{2}$ |  |  | 146.4 | 131.6 | 192.0 | 199.0 | 174.0 |
| $\mathrm{~B}_{1}$ |  |  |  | 80.8 | 158.7 | 108.4 | 174.1 |
| $\mathrm{~A}_{1}$ |  |  |  | 264.8 | 234.3 | 282.7 |  |
| $\mathrm{~B}_{2}$ |  |  |  | 87.6 | 41.7 | 66.0 |  |
| $\mathrm{~B}_{1}$ |  |  |  | 69.3 | 9.5 | 66.9 |  |

${ }^{a}$ CCSD/aug-cc-DVP ref. 2.

The calculated IR spectra for all species are marked by two strong NN stretches at $\sim 1300$ (asym, $B_{2}$ ) and $\sim 1250 \mathrm{~cm}^{-1}$ (sym, $A_{1}$ ) and a prominent ring-metal atom stretch $\left(A_{1}\right)$ between 600 and $200 \mathrm{~cm}^{-1}$. Results for the B3LYP $/ / 6-311+G(d)$ method are given in Table 2. The symmetric and asymmetric stretches are similar within element groups $1(\mathrm{Na}, \mathrm{K}), 2(\mathrm{Mg}, \mathrm{Ca})$, and 12 $(\mathrm{Zn})$. The values calculated at the RB3LYP//6-31G(d) level are generally $10 \mathrm{~cm}^{-1}$ higher. Results for the isolated pentazole anion are similar to those reported by Bartlett et al. ${ }^{2}$ calculated with the coupled cluster CCSD method with an augmented DVP basis set.
In order to calibrate the theoretical method the N atom chemical shifts were calculated for 1- and 2-methyltetrazole and compared to experiment. The results are given in Table 4. A linear regression analysis on the eight data points gives the formula $\sigma_{\text {exp }}=\left(1.0683 \sigma_{\text {calc }}-4.70\right)$ with a correlation coefficient of 0.9988 . The value of the intercept -4.70 and a slope near 1.0 indicate that the B3LYP/6-311++G(2d,p) GIAO calculation on a B3LYP $/ 6-31 \mathrm{G}(\mathrm{d})$ optimised structure gives values very close to experimental ones and much closer than those using the CHF method $\left(0.8804 \sigma_{\text {calc }}+112.56\right) .{ }^{9}$



The absolute shielding value for the N atom in nitromethane is calculated to be -158.225 ppm using the B3LYP//6$311++\mathrm{G}(2 \mathrm{~d}, \mathrm{p})$ method on the B3LYP//6-31G(d) optimised structure. The theoretical value for the N atom in $\mathrm{NH}_{3}$ is +258.440 ppm . Thus, 100.215 ppm needs to be added to the above values for comparison to external $\mathrm{NH}_{3}$ rather than nitromethane. The reported isotropic chemical shifts in ppm were calculated as -158.225 minus the absolute shielding value calculated for each N atom in the molecule in question. There are three distinct N atoms in the II and IIa structures (also in I for Zn ). The values in ppm for the N atom chemical shifts are reported in Table 3, using the order of closest to furthest from the metal atom to number the N atoms. As seen in Table 1, the energy differences between the I and II forms vary from 4 to 9 $\mathrm{kcal} \mathrm{mol}{ }^{-1}$. Although this barrier might prevent intramolecular scrambling, intermolecular exchange of the metal atom in solution is expected even at low temperatures. Thus, only one NMR ${ }^{14} \mathrm{~N}$ or ${ }^{15} \mathrm{~N}$ signal is expected in solution.

It can be seen in Table 3 that there is very little difference in chemical shift for the $\operatorname{Mg}\left(\mathrm{N}_{5}\right)_{2}$ and $\mathrm{MgClN}_{5}$ (II) species. The differences in values for atoms within the ring of the various species reflect the charge distribution. N atoms bonded to
the metal have more charge and thus the chemical shifts are more negative. This effect is seen to increase for the metals in the order of $\mathrm{K}<\mathrm{Na} \sim \mathrm{Ca}<\mathrm{Mg}<\mathrm{Zn}$. The much larger negative shift in Zn is most likely due to the $\mathrm{Zn}-\mathrm{N}$ bond in conformation I.
There have been structural and mechanistic studies ${ }^{3}$ involving ${ }^{15} \mathrm{~N}$ labelling and the addition of the azide ion to the aryl diazonium ion (Scheme 4) and the recovery of the aryl azide.


Scheme 4 Numbering for the ${ }^{15} \mathrm{~N}$ atom experimental and theoretical chemical shifts.

Each study proposed a pentazole intermediate as the origin of a transitory peak in the ${ }^{15} \mathrm{~N}$ spectrum. The N atom shifts have been calculated for all species in Scheme 4 and their values are reported in Table 3. For the $p$-methoxyphenyl series (c) the calculations indicated two rotational isomers due to restricted rotation of the aryl-oxygen bond and the mean values of the ${ }^{15} \mathrm{~N}$ shifts are given in Table 3. These rotational isomers were not seen in either the proton or ${ }^{15} \mathrm{~N}$ spectra. The calculations indicated rotational barriers of the order of $3 \mathrm{kcal} \mathrm{mol}^{-1}$ and hence rapid exchange between the forms would be expected. ${ }^{15} \mathrm{~N}$-Labelled samples of the compounds $\mathbf{1 c}, 3 \mathrm{c}$ and $\mathbf{4 c}$ were obtained by diazotising $p$-methoxyaniline with $\mathrm{Na}^{15} \mathrm{NO}_{2}$ thereby putting a label at N 2 . The other ${ }^{15} \mathrm{~N}$ shifts listed in Table 3 for these compounds were measured in natural abundance. The calculated ${ }^{15} \mathrm{~N}$ shifts are in good agreement with the measured values except for N 2 of the diazonium ion 1c (Table 3). In this case the presence of the chloride anion in the solution is expected to increase the electron density and hence the shielding at N2. This may also account for the small increased shielding measured at N1 for the species 1c (Table 3). The influence of the counter ion is also evident in the calculated shifts for the $\mathrm{N}_{5}$ anions (Table 3).

The principal excitation energies ( nm ) and oscillator strengths (au) calculated with the TD-DFT, B3LYP//6-311+ G(d) method are given in Table 5. State symmetries are given

Table 3 Theoretical (and experimental) ${ }^{a}$ values of N atom shifts relative to external nitromethane (ppm)

| Molecule | N1 | N2 | N3 | N4 | N5 | $\mathrm{NMe}_{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1a | $\begin{gathered} -147.4 \\ (-148.3) \end{gathered}$ | -24.7 | - | - | - | - |
| 1b | -132.3 | -11.3 | - | - | - | -277.9 |
| 1c | $\begin{aligned} & -142.8 \\ & (-) \end{aligned}$ | $\begin{gathered} -12.6 \\ (-61.2)^{f} \end{gathered}$ | - | - | - | - |
|  | $(-147.1)^{b}$ | $(-59.7)^{\text {b }}$ |  |  |  |  |
| 2a | 66.5 | 88.9 | -210.9 | -144.5 | -94.0 |  |
| 3a | $\begin{gathered} -85.3 \\ (-82.7) \end{gathered}$ | -27.3 | 18.1 | - | - |  |
| 3b | $\begin{gathered} -83.1 \\ (-80.0) \end{gathered}$ | $\begin{gathered} -31.3 \\ (-27.1) \end{gathered}$ | $\begin{aligned} & 14.0 \\ & (4.9) \end{aligned}$ | - | - | $\begin{array}{r} -339.9 \\ (-3246 \end{array}$ |
| 3c | -84.1 ${ }^{\text {c }}$ | -29.2 ${ }^{\text {c }}$ | $-18.8^{\text {c }}$ |  |  |  |
|  | (-) | $(-28.1)^{f}$ | (-) |  |  |  |
| 4a | $\begin{gathered} -307.4 \\ (-288.1) \end{gathered}$ | -146.2 | -143.7 | - | - | - |
| 4b | $\begin{gathered} -312.4 \\ (-292.2) \end{gathered}$ | $\begin{gathered} -141.6 \\ (-134.6) \end{gathered}$ | $\begin{gathered} -142.9 \\ (-146.2) \end{gathered}$ | - | - | $\begin{gathered} -346.5 \\ (-333.2) \end{gathered}$ |
| 4c | $\begin{aligned} & -310.4^{c} \\ & (-293.0)^{d} \\ & (-) \end{aligned}$ | $\begin{aligned} & -144.2^{c} \\ & (-137.3)^{d f} \\ & (-139.7) \end{aligned}$ | $\begin{aligned} & -144.3^{c} \\ & (-149.3)^{d} \\ & (-) \end{aligned}$ |  |  |  |
| $\mathrm{N}_{5}$ Anion | -1.7 | -1.7 | -1.7 | -1.7 | -1.7 | - |
| $\mathrm{NaN}_{5}$ (II) | -28.7 | 3.0 | 14.9 | - | - 17.6 | - |
| $\mathrm{NaN}_{5}$ (III) | 17.6 | 17.6 | 17.6 | 17.6 | 17.6 |  |
| $\mathrm{KN}_{5}$ | -18.8 | 2.2 | 15.5 | - | - | - |
| $\mathrm{Mg}\left(\mathrm{N}_{5}\right)_{2}$ | -48.2 | 6.3 5 | 24.9 | - | - | - |
| $\mathrm{MgN}_{5} \mathrm{Cl}$ | -47.4 | 5.8 | 23.9 | - | - | - |
| $\mathrm{CaN}_{5} \mathrm{Cl}$ | $-27.2$ | 6.3 | $29.4$ | - | - | - |
| $\mathrm{ZnN}_{5} \mathrm{Cl}$ (I) | -96.1 ${ }^{\text {e }}$ | -7.9 | 14.5 | - | - | - |

${ }^{a}$ From ref. 3 or measured herein as indicated. ${ }^{b}$ Measured herein in $\mathrm{CD}_{3} \mathrm{OD}-\mathrm{D}_{2} \mathrm{O}(70: 30 \mathrm{v} / \mathrm{v})$ for natural abundance ${ }^{15} \mathrm{~N}$ at $0{ }^{\circ} \mathrm{C} .{ }^{c}$ Values $\pm 0.3$. ${ }^{d}$ Measured herein in $\mathrm{CDCl}_{3}$ for natural abundance ${ }^{15} \mathrm{~N}$ at probe temperature. ${ }^{e} \mathrm{~N} 1$ in contact with $\mathrm{Zn} .{ }^{f}$ Measured herein in $\mathrm{CD}_{3} \mathrm{OD}$ for ${ }^{15} \mathrm{~N}$ labelled samples at $-35^{\circ} \mathrm{C}$.

Table 4 Theoretical and experimental ${ }^{9}$ values ( ppm ) of N atom chemical shifts for 1-methyltetrazole and 2-methyltetrazole relative to nitromethane

| Substituent | Atom | DFT | Experiment |
| :---: | :--- | ---: | :--- |
|  |  |  |  |
| 1-Methyl | 1 | -163.1 | -159.6 |
|  | 2 | -6.3 | -8.4 |
|  | 3 | 34.2 | 23.3 |
| 2-Methyl | 4 | -41.9 | -43.2 |
|  | 1 | -78.4 | -73.1 |
|  | 2 | -107.7 | -107.0 |
|  | 3 | 9.5 | 3.4 |
|  | 4 | -42.3 | -41.3 |

for the anion based on the $D_{5 \mathrm{~h}}$ group, as well as those based on the equivalent $C_{2 \mathrm{v}}$ symmetries when the cations are present. Those states reported in Table 5 are of three types. The first group of six states is where the principal excitation is to the metal s AO from the four highest lone pair MOs ${ }^{1}\left(\mathrm{~ns}^{*}\right)$ and the two highest $\pi \operatorname{MOs}^{1}\left(\pi \mathrm{~s}^{*}\right)$. The next is a group of four ${ }^{1}\left(\mathrm{n} \pi^{*}\right)$ MOs and the last is a group of four ${ }^{1}\left(\pi \pi^{*}\right)$ MOs where these named excitations are the principal excitation within that state. A 77 state TD-DFT calculation had to be done to assure that the above 14 states were all found. Other states included excitations such as from the Cl atom lone pairs or to the metal p and d virtual AOs.

The ${ }^{1}\left(\pi \pi^{*}\right)$ transitions in the free anion are similar to those in benzene in that the lowest lying ones both involve excitations from a degenerate pair of $\pi \mathrm{MOs}$ to another degenerate pair. In benzene the $D_{6 h}$ symmetry gives an occupied $\mathrm{e}_{1 \mathrm{~g}}$ pair and an unoccupied $e_{2 u}$ pair. The direct product of these transitions leads to three states, $B_{1 u}+B_{2 u}+E_{1 u}$. In free pentazole anion the $D_{5 \mathrm{~h}}$ symmetry leads to a $\pi$ occupied $\mathrm{e}_{1}{ }^{\prime \prime}$ pair and a $\pi$ unoccupied $\mathrm{e}_{2}{ }^{\prime \prime}$ pair.

The direct product in this case is $\mathrm{e}_{1}{ }^{\prime \prime} \times \mathrm{e}_{2}{ }^{\prime \prime}=\mathrm{E}_{2}{ }^{\prime}+\mathrm{E}_{1}{ }^{\prime}$. Similarly for the four highest ${ }^{1}\left(n \pi^{*}\right)$ transitions the lone pair HOMOs are an $\mathrm{e}_{1}{ }^{\prime}$ pair. This gives the direct products:
$\mathrm{e}_{1}{ }^{\prime} \times \mathrm{e}_{2}{ }^{\prime \prime}=\mathrm{E}_{2}{ }^{\prime \prime}+\mathrm{E}_{1}{ }^{\prime \prime}$. Upon distortion of one atom in the free anion to give $C_{2 \mathrm{v}}$ symmetry or in the presence of a metal cation in the I or II conformation ( $c f$. Fig. 1), the orbital degeneracies are lifted. Orbitals with $\mathrm{e}_{1}{ }^{\prime \prime}$ or $\mathrm{e}_{2}{ }^{\prime \prime}$ symmetries lead to $\mathrm{a}_{2}$ or $\mathrm{b}_{1}$ symmetries; the $\mathrm{e}_{1}{ }^{\prime}$ symmetry of the lone pair HOMOs lead to two lone pair MOs of $a_{1}$ and $b_{2}$ symmetries.

The ${ }^{1}\left(\mathrm{n} \pi^{*}\right)$ transitions reported here are electronically forbidden as the $\mathrm{E}_{2}{ }^{\prime \prime}$ and $\mathrm{E}_{1}{ }^{\prime \prime}$ states do not transform like an in-plane axis. This is also the case for the ${ }^{1}\left(\pi \pi^{*}\right) \mathrm{E}_{2}{ }^{\prime}$ state. The ${ }^{1}\left(\pi \pi^{*}\right) \mathrm{E}_{1}{ }^{\prime}$ state transforms like $(x, y)$ and in consequence this transition is polarized in the plane of the anion, giving an oscillator strength of 0.373 au . In the presence of the cation, however, the $C_{2 v}$ group lifts the symmetry restraints on all but the new $\mathrm{A}_{2}$ states.

The ${ }^{1}\left(\mathrm{n} \pi^{*}\right)$ transition energies vary significantly according to the metal counterion. This can also be seen at the level of the DFT orbital energies, Table 6. The orbital energy differences for $\Delta \mathrm{ns}^{*}, \Delta \mathrm{n} \pi^{*}$, and $\Delta \pi \pi^{*}$ correspond to the excitation energies calculated with the TD-DFT method. The change in the anion orbital levels induced by the presence of the metal counterion is also significant. The amount of lowering of the lone pair and $\pi$ HOMOs follows the order of $\mathrm{K}<\mathrm{Na} \ll$ $\mathrm{Ca}<\mathrm{Zn}<\mathrm{Mg}$. Group I metals are more electropositive than group II, and K more than Na . Since Ca lies below Mg , it can be expected to interact less with the ring. The Zn atom is in the $\mathbf{I}$ conformation where the $\mathbf{M}-\mathbf{N}$ distance is shorter than in the II conformation and thus the interaction is expected to be greater, giving the anomalous position for Zn in the series.

In summary, it is seen that the nature of the cation has an influence on the electronic, vibrational, and NMR spectra of the pentazole anion. A complex with the metal ion in the plane of the ring is energetically favoured. This induces an alternating bond length structure of $C_{2 \mathrm{v}}$ symmetry and redistribution of charge in the anion. Further theoretical studies are in progress on the effect of solvent and on the pericyclic decomposition of the pentazole ring in the presence of a cation.

Table 5 Principal excitation energies ( nm ) and oscillator strengths (au) calculated with the B3LYP//6-311+G(d) method. State symmetries are given for the anion based on the $D_{5 \mathrm{~h}}$ group; those based on the equivalent $C_{2 \mathrm{v}}$ symmetries are given in parentheses

|  | $\mathrm{N}_{5}{ }^{-}$ | NaN 5 | $\mathrm{KN}_{5}$ | $\mathrm{MgN}_{5} \mathrm{Cl}$ | $\mathrm{CaN}_{5} \mathrm{C}$ |  | $\mathrm{ZnN}_{5} \mathrm{Cl}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{1}\left(\mathrm{n}, \mathrm{s}^{*}\right)-\left(\mathrm{A}_{1}\right)$ | - | 255.90 .038 | 251.00 .029 | 193.60 .046 | 244.4 | 0.006 | 211.0 | 0.040 |
| ${ }^{1}\left(\mathrm{n}, \mathrm{s}^{*}\right)-\left(\mathrm{B}_{2}\right)$ | - | 255.50 .001 | 249.90 .001 | 191.80 .002 | 223.8 | 0.003 | 209.0 | 0.02 |
| ${ }^{1}\left(\pi, \mathrm{~s}^{*}\right)-\left(\mathrm{B}_{1}\right)$ | - | 243.1 0.0.003 | 235.10 .002 | 180.30 .009 | 208.8 | 0.005 | 208.7 | 0.014 |
| ${ }^{1}\left(\mathrm{n}, \mathrm{s}^{*}\right)$ - $\left(\mathrm{A}_{1}\right)$ | - | 240.00 .045 | 236.80 .023 | 182.80 .045 | 224.2 | 0.000 | 170.2 | 0.004 |
| ${ }^{1}\left(\pi, \mathrm{~s}^{*}\right)-\left(\mathrm{A}_{2}\right)$ | - | 237.00 .0 | 231.80 .0 | 178.00 .0 | 205.4 | 0.0 | 191.2 | 0.0 |
| ${ }^{1}\left(\mathrm{n}, \mathrm{s}^{*}\right)$ - $\left(\mathrm{B}_{2}\right)$ | - | 208.50 .008 | 208.40 .003 | 172.00 .028 | 174.9 | 0.001 | 194.6 | 0.005 |
| ${ }^{1}\left(\mathrm{n}, \pi^{*}\right) \mathrm{E}_{2}{ }^{\prime \prime}\left(\mathrm{A}_{2}\right)^{a}$ | 188.40 .0 | 197.90 .0 | 197.20 .0 | 205.50 .0 | 213.7 | 0.0 | 183.4 | 0.0 |
| ${ }^{1}\left(\mathrm{n}, \pi^{*}\right) \mathrm{E}_{2}{ }^{\prime \prime}\left(\mathrm{B}_{1}\right)^{b}$ | 188.40 .0 | 197.70 .004 | 197.10 .003 | 204.00 .005 | 209.4 | 0.004 | 218.9 | 0.004 |
| ${ }^{1}\left(\mathrm{n}, \pi^{*}\right) \mathrm{E}_{1}{ }^{\prime \prime}\left(\mathrm{B}_{1}\right)^{c}$ | 186.70 .0 | 188.50 .001 | 184.10 .000 | 185.30 .008 | 192.40 | 0.001 | 176.80 | 0.007 |
| ${ }^{1}\left(\mathrm{n}, \pi^{*}\right) \mathrm{E}_{1}{ }^{\prime \prime}\left(\mathrm{A}_{2}\right)^{d}$ | 186.70 .0 | 187.90 .0 | 182.90 .0 | 184.60 .0 | 189.0 | 0.0 | 215.6 | 0.0 |
| ${ }^{1}\left(\pi, \pi^{*}\right) \mathrm{E}_{2}^{\prime}\left(\mathrm{B}_{2}\right)^{e}$ | 166.00 .0 | 166.50 .021 | 168.80 .014 | 172.00 .028 | 180.4 | 0.028 | 170.7 | 0.13 |
| ${ }^{1}\left(\pi, \pi^{*}\right) \mathrm{E}_{2}{ }^{\prime}\left(\mathrm{A}_{1}\right)^{f}$ | 166.00 .0 | 163.50 .042 | 164.00 .017 | 160.50 .009 | 171.5 | 0.041 | 173.9 | 0.169 |
| ${ }^{1}\left(\pi, \pi^{*}\right) \mathrm{E}_{1}{ }^{\prime}\left(\mathrm{A}_{1}\right)^{g}$ | 153.20 .353 | 152.30 .088 | 147.70 .270 | 153.10 .467 | $159.0^{i}$ | 0.197 | 186.7 | 0.086 |
| ${ }^{1}\left(\pi, \pi^{*}\right) \mathrm{E}_{1}{ }^{\prime}\left(\mathrm{B}_{2}\right)^{h}$ | 153.20 .353 | 150.30 .023 | 147.50 .185 | 152.50 .045 | $161.0^{i}$ | 0.085 | 147.6 | 0.462 |

${ }^{a}$ Major excitation: occupied $\mathrm{A}_{1}$ to virtual $\mathrm{A}_{2} \cdot{ }^{b}$ Major excitation: occupied $\mathrm{B}_{2}$ to virtual $\mathrm{A}_{2} \cdot{ }^{c}$ Major excitation: occupied $\mathrm{A}_{1}$ to virtual $\mathrm{B}_{1} .{ }^{d}$ Major excitation: occupied $\mathrm{B}_{2}$ to virtual $\mathrm{B}_{1} .{ }^{e}$ Major excitation: occupied $\mathrm{B}_{1}$ to virtual $\mathrm{A}_{2} \cdot{ }^{f}$ Major excitation: occupied $\mathrm{A}_{2}$ to virtual $\mathrm{A}_{2} .{ }^{g}$ Major excitation: occupied $\mathrm{B}_{1}$ to virtual $\mathrm{B}_{1} \cdot{ }^{h}$ Major excitation: occupied $\mathrm{A}_{2}$ to virtual $\mathrm{B}_{1} \cdot{ }^{i}$ Middle state of three states containing the major excitation.

Table 6 Orbital energies and energy differences calculated with the B3LYP//6-311+G(d) method

|  | Anion | $\mathrm{NaN}_{5}$ |  | $\mathrm{KN}_{5}$ | $\mathrm{MgClN}_{5}$ | $\mathrm{CaClN}_{5}$ | $\mathrm{ZnClN}_{5}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| S | - | -0.0839 | -0.0683 | -0.0845 | -0.1033 | -0.0942 |  |
| $\mathrm{HO}-\mathrm{n}$ | -0.1148 | -0.2827 | -0.2692 | -0.3448 | -0.3228 | -0.3378 |  |
| $\mathrm{HO}-\pi$ | -0.1247 | -0.2981 | -0.2854 | -0.3701 | -0.3522 | -0.3544 |  |
| $\mathrm{LU} \pi^{*}$ | +0.1864 | +0.0107 | +0.0341 | -0.0719 | -0.0670 | -0.0822 |  |
| $\Delta \mathrm{nS}$ | - | 0.1988 | 0.2009 | 0.2603 | 0.2195 | 0.2436 |  |
| $\Delta \pi \mathrm{~S}$ | - | 0.2142 | 0.2171 | 0.2856 | 0.2489 | 0.2602 |  |
| $\Delta \mathrm{n} \pi^{*}$ | 0.3012 | 0.2934 | 0.3033 | 0.2729 | 0.2558 | 0.2556 |  |
| $\Delta \pi \pi^{*}$ | 0.3111 | 0.3088 | 0.3195 | 0.2982 | 0.2852 | 0.2722 |  |

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## References

1 R. Huisgen and I. Ugi, Chem. Ber., 1957, 90, 2914; R. Huisgen and I. Ugi, Chem. Ber., 1958, 91, 531; for reviews see (a) I. Ugi Comprehensive Heterocyclic Chemistry, Series eds. A. R. Katritzky and C. W. Rees, Pergamon Press, Oxford, 1984, vol. 5, p. 839; (b) R. N. Butler, Comprehensive Heterocyclic Chemistry II, Series eds., A. R. Katritzky, C. W. Rees and E. F. V. Scriven, Pergamon Press, Oxford, 1996, vol. 4 (ed., R. C. Storr), p. 897.
2 M. T. Nguyen, M. Sana, G. Leroy and J. Elguero, Can. J. Chem., 1983, 61, 1435; M. T. Nguyen, M. A. McGinn, A. F. Hegarty and J. Elguero, Polyhedron, 1985, 4, 1721; R. J. Bartlett, Chem. Ind. (London), 2000, 140; K. F. Ferris and R. J. Bartlett, J. Am. Chem. Soc., 1992, 114, 8302
3 R. N. Butler, A. Fox, S. Collier and L. A. Burke, J. Chem. Soc., Perkin Trans. 2, 1998, 2243; R. Müller, J. D. Wallis and W. von Philipsborn, Angew. Chem., Int. Ed. Engl., 1985, 24, 513.
4 Gaussian98, Revision A.7, M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, V. G Zakrzewski, J. A. Montgomery, Jr., R. E. Stratmann, J. C. Burant,
S. Dapprich, J. M. Millam, A. D. Daniels, K. N. Kudin, M. C. Strain, O. Farkas, J. Tomasi, V. Barone, M. Cossi, R. Cammi, B. Mennucci, C. Pomelli, C. Adamo, S. Clifford, J. Ochterski, G. A. Petersson, P. Y. Ayala, Q. Cui, K. Morokuma, D. K. Malick, A. D. Rabuck, K. Raghavachari, J. B. Foresman, J. Cioslowski, J. V. Ortiz, A. G. Baboul, B. B. Stefanov, G. Liu, A. Liashenko, P. Piskorz, I. Komaromi, R. Gomperts, R. L. Martin, D. J. Fox, T. Keith, M. A. Al-Laham, C. Y. Peng, A. Nanayakkara, C. Gonzalez, M. Challacombe, P. M. W. Gill, B. Johnson, W. Chen, M. W. Wong, J. L. Andres, C. Gonzalez, M. Head-Gordon, E. S. Replogle and J. A. Pople, Gaussian, Inc., Pittsburgh PA, 1998.

5 A. D. Becke, Phys. Rev. A, 1988, 38, 3098; C. Lee, W. Yang and R. G. Parr, Phys. Rev. B, 1988, 37, 785.
6 R. Seeger and J. A. Pople, J. Chem. Phys., 1977, 66, 3045; R. Bauernschmitt and R. Aldrichs, J. Chem. Phys., 1996, 104, 9047.

7 S. Portmann and H. P. Lüthi, MOLEKEL: An Interactive Molecular Graphics Tool, Chimia, 2000, 54, 766
8 K. Wolinski, J. F. Hilton and P. Pulay, J. Am. Chem. Soc., 1990, 112, 8251; J. L. Dodds, R. McWeeny and A. J. Sadlej, Mol. Phys., 1980, 41, 1419; R. Ditchfield, Mol. Phys., 1974, 27, 789.
9 M. Witanowski, Z. Biedrzycka, W. Sicinska and Z. Grabowski, J. Magn. Reson., 1998, 131, 54.

10 R. E. Stratmann, G. E. Scuseria and M. J. Frisch, J. Chem. Phys., 1998, 109, 8218; R. Bauernschmitt and R. Aldrichs, Chem. Phys. Lett., 1996, 256, 454; M. E. Casida, C. Jamorski, K. C. Casida and D. R. Salahub, J. Chem. Phys. , 1998, 108, 4439.

